

# A Tale of Two Mergers

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> Contrast evolutions of two distinct mergers:

(1) A black hole & a normal star (NS)

(2) A black hole & a self-bound star (SBS)

$$\alpha = \frac{d \ln R}{d \ln M} \quad \left\{ \begin{array}{l} \leq 0 \text{ for a NS} \\ \geq 0 \text{ for a SBS (SQM?)} \end{array} \right.$$

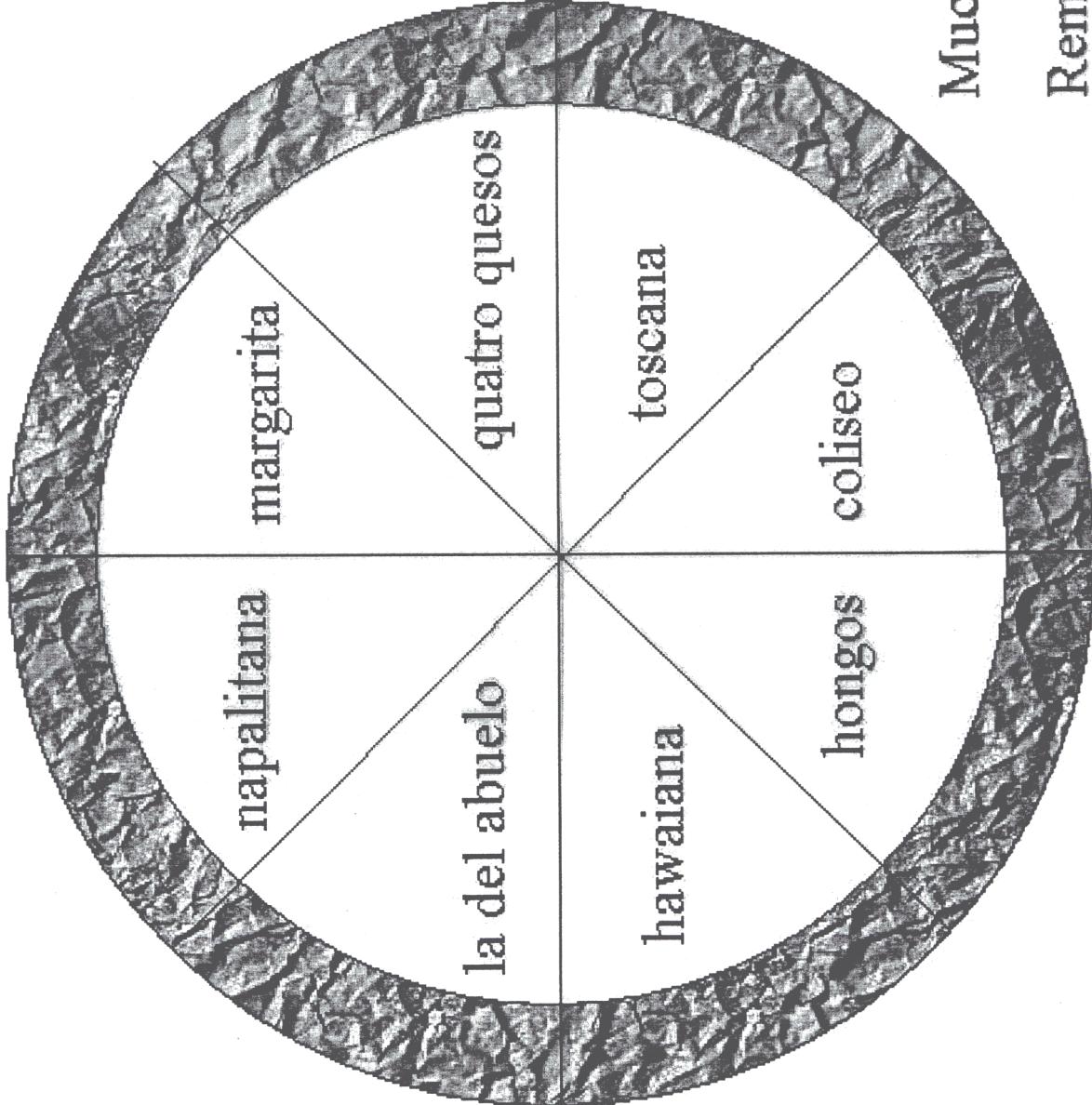
$\nearrow$   
"EOS"

Explore astrophysical consequences  
of differences in  $\alpha$  in

(1) Time scales

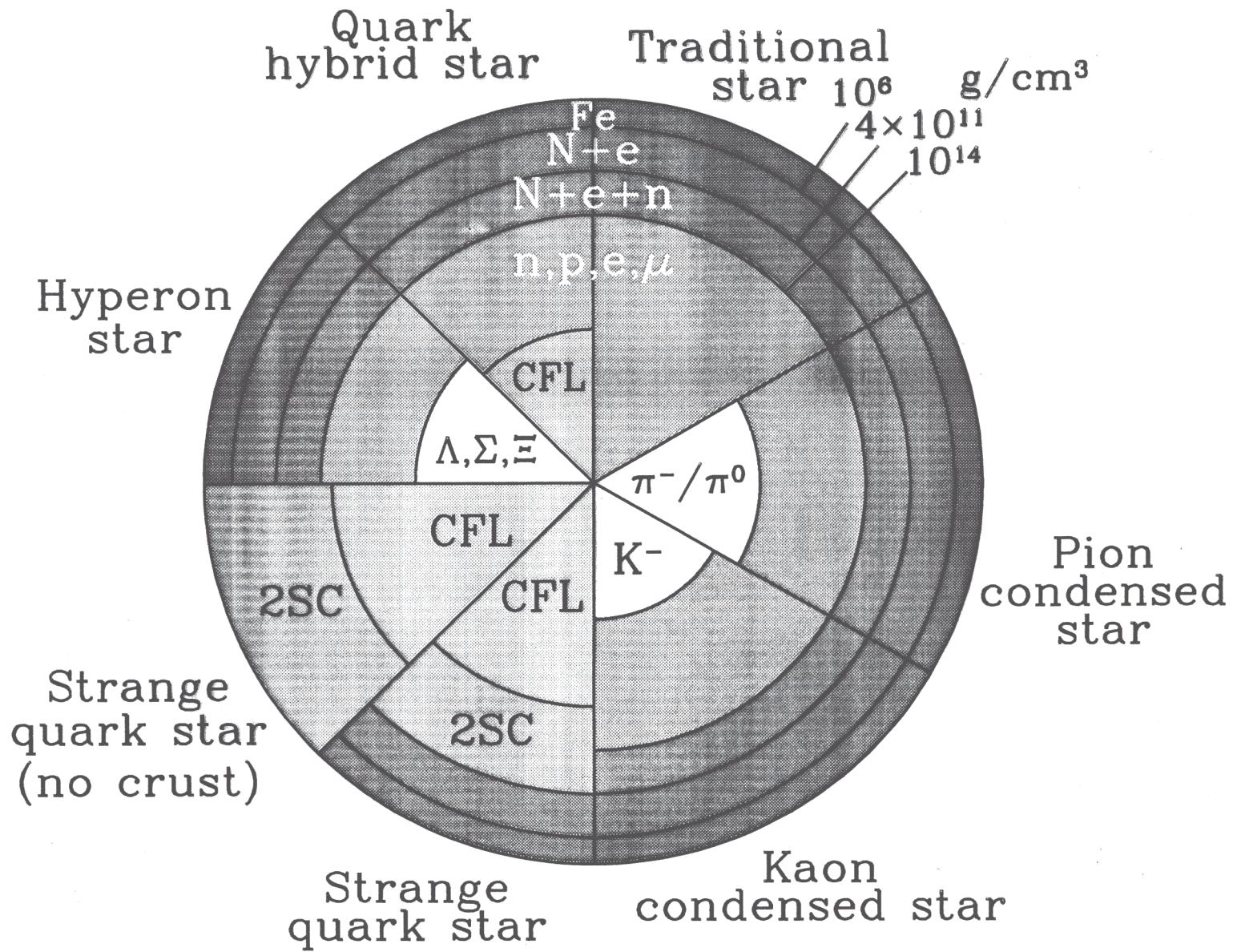
(2) Gravity wave signals

# de LENA

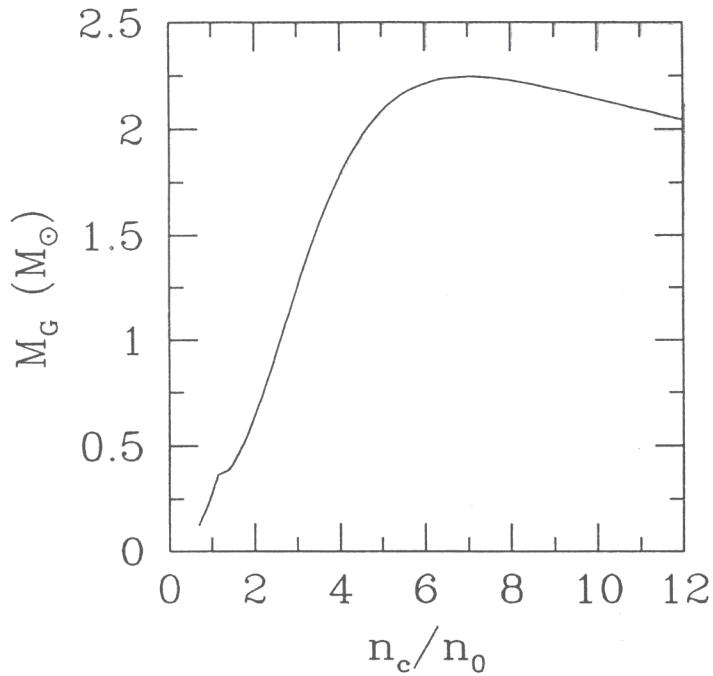


Muchas Gracias

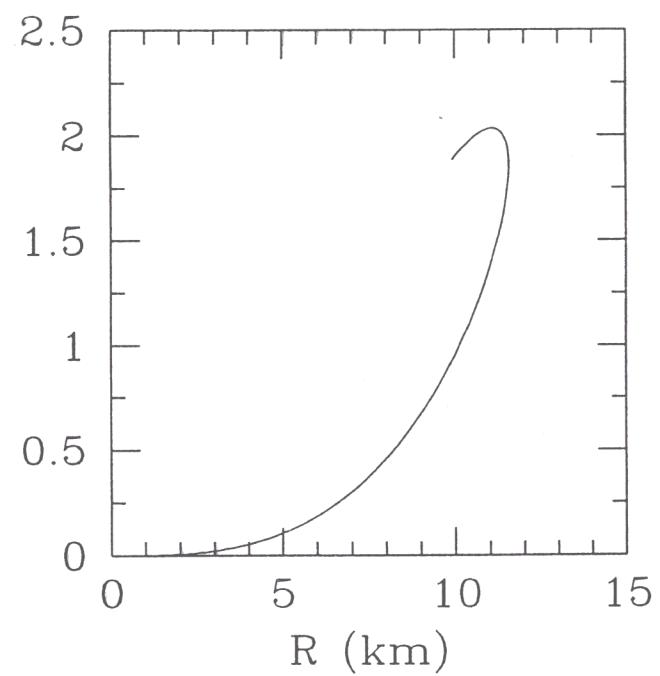
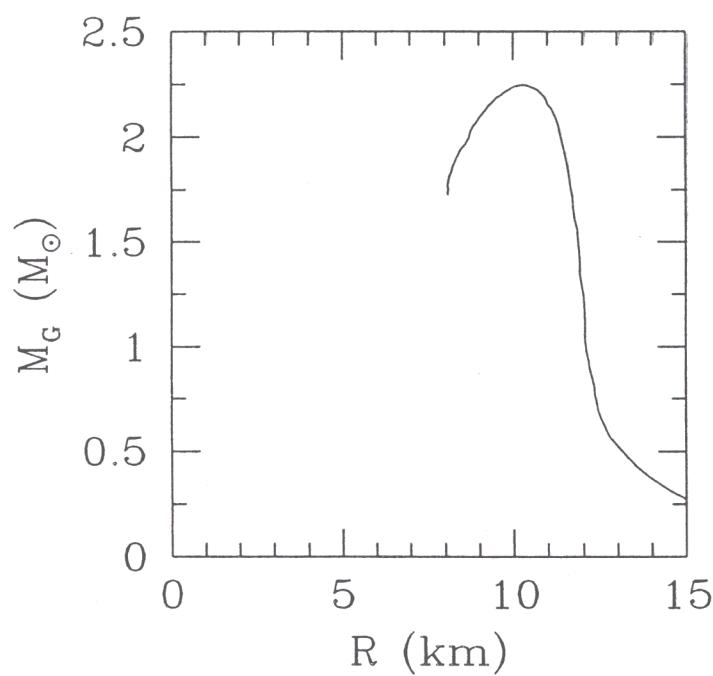
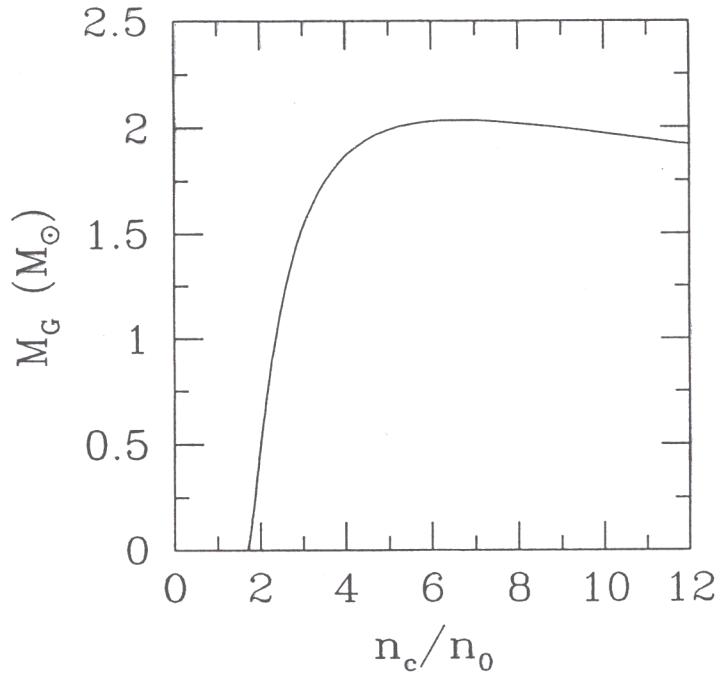
Remy Avila!

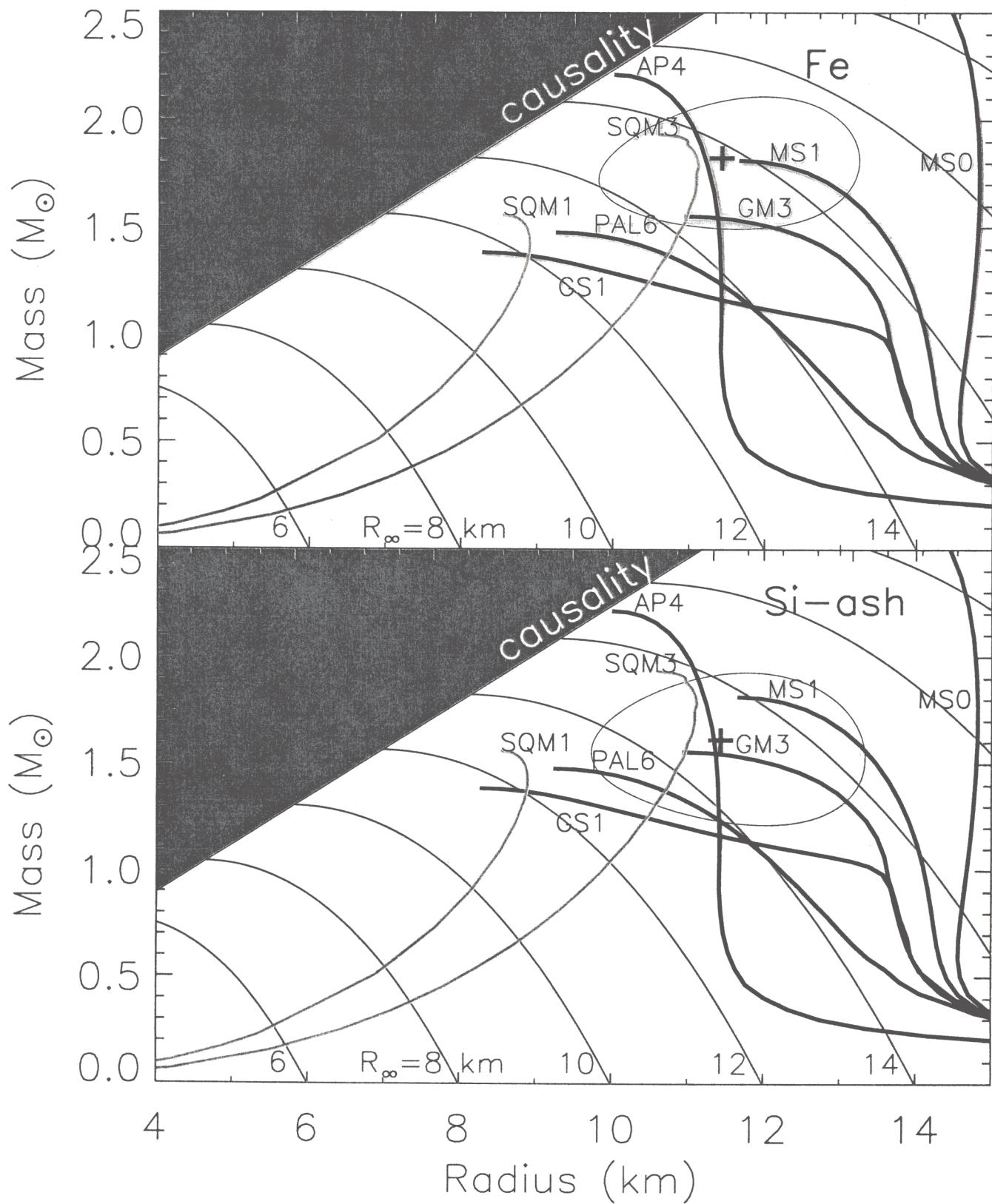


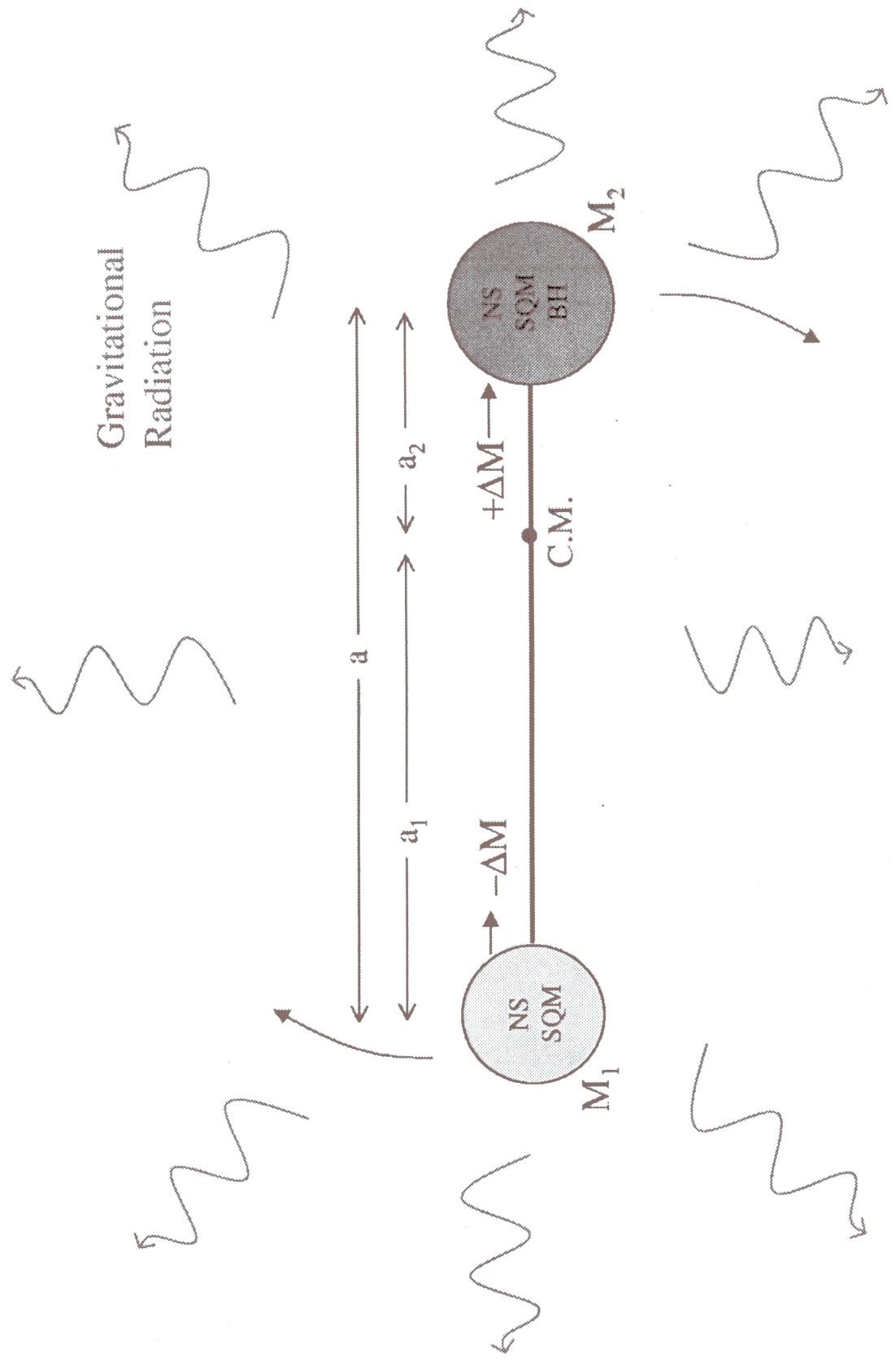
Normal Star



Self-Bound Star







$$L_{GW} = \frac{1}{5} \frac{G}{c^5} \langle \ddot{I}_{jk} \ddot{I}_{jk} \rangle$$

$$M = M_1 + M_2$$

$$\mu = \frac{M_1 M_2}{M}$$

$$= \frac{32}{5} \frac{G^4}{c^5} \frac{M^3 \mu^2}{a^5}$$

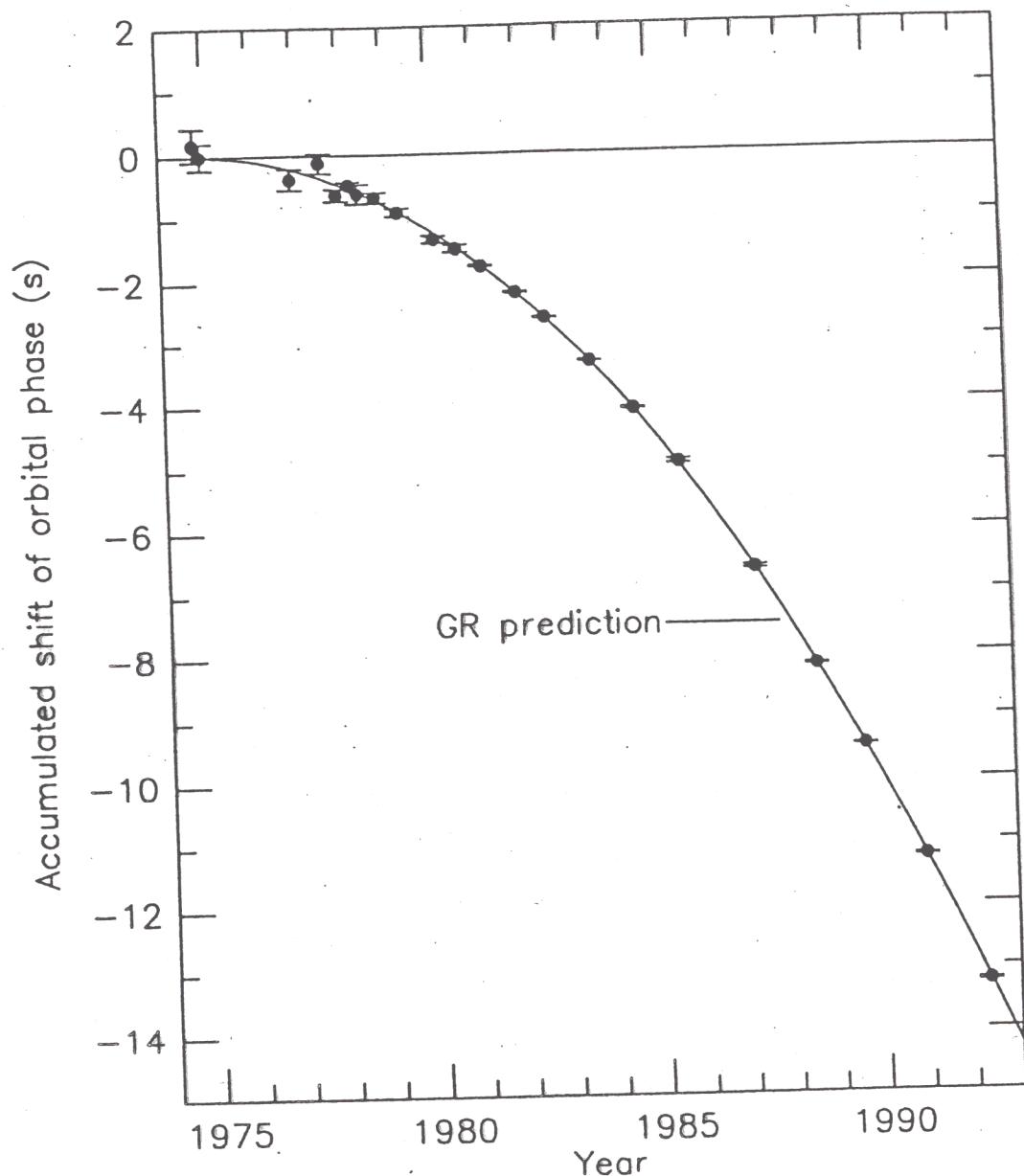
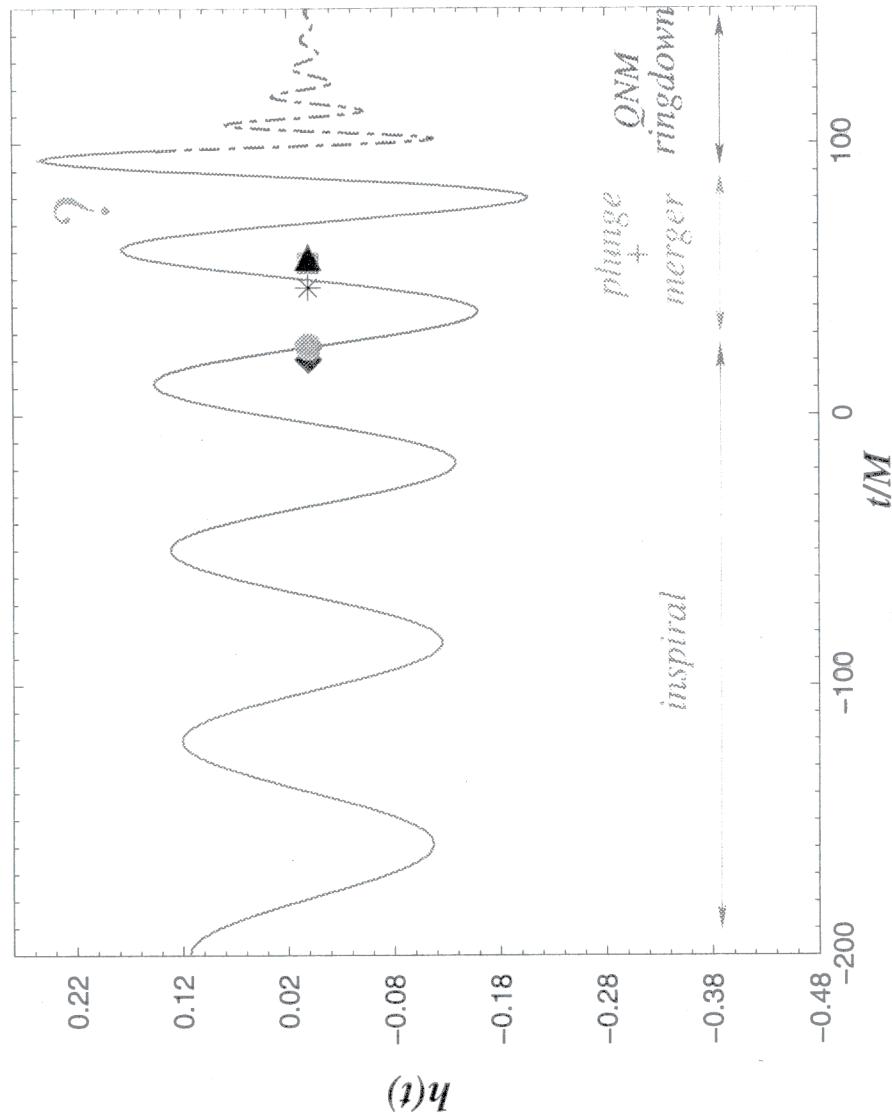


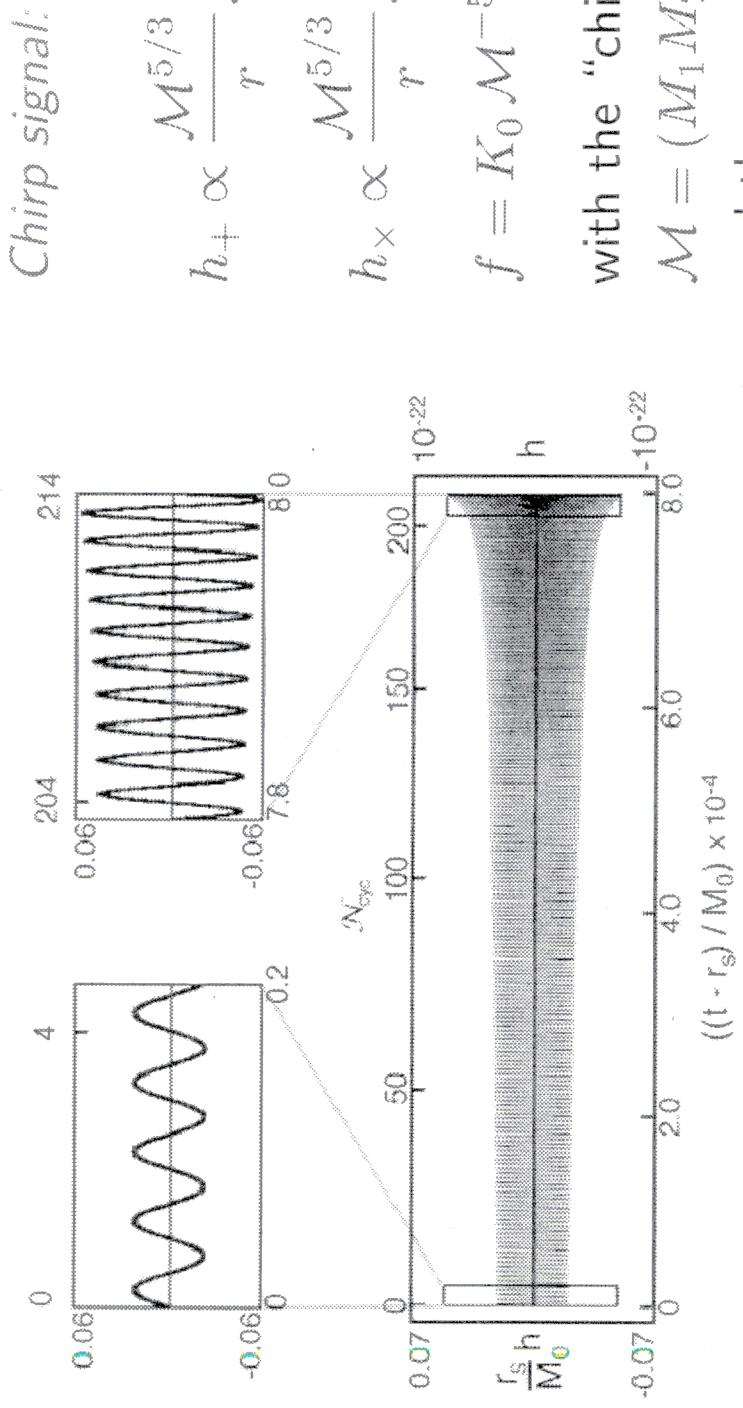
FIG. 10. Accumulated shift of the times of periastron in the PSR 1913+16 system, relative to an assumed orbit with constant period. The parabolic curve represents the general relativistic prediction for energy losses from gravitational radiation.

## Gravitational waveform



[from Buonanno & Damour, PRD 62, 064015 (2000) ]

## Inspiral waveform



with the "chirp mass":

$$\mathcal{M} = (M_1 M_2)^{3/5} (M_1 + M_2)^{-1/5}$$

and the constant:

$$K_0 = \frac{5^{3/8}}{8\pi} \left(\frac{c^3}{G}\right)^{5/8}$$

[Duez, Baumgarte & Shapiro, PRD 63, 084030 (2001)]

- Evolution of a compact binary

Mass ratio:  $q = M_2 / M_1$

$$M_1 = M_{\text{NS}} \quad \text{SQM}$$

$$M_2 = M_{\text{BH}}$$

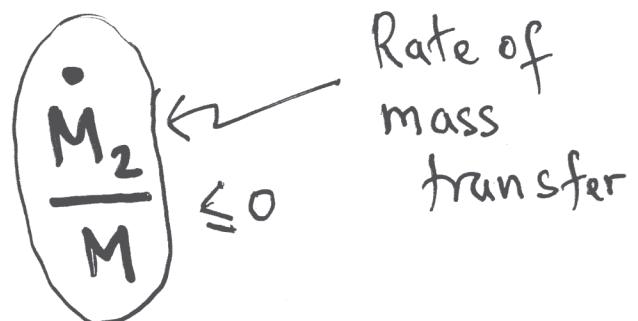
Total mass:  $M = M_1 + M_2$

Reduced mass:  $\mu = M_1 M_2 / M$

> Orbital angular momentum:

$$J^2 = G \mu^2 M a$$

$$\dot{J} = \frac{G^{1/2} M_2^{3/2} a^{1/2}}{q^{1/2} (1+q)^{1/2}}$$



$$* \left[ (1-q) + \underbrace{\frac{1}{2} \frac{d \ln a}{d \ln M_2}}_{\nearrow} \right]$$

$\nearrow$   
change of orbit separation

To conserve angular momentum

EOS

$$\dot{J} = \frac{G^{1/2} a^{1/2} M^{3/2}}{(1+q)^3} \dot{q} \left[ \frac{\alpha}{2} + \frac{5}{6} - q \right]$$

$$\gg \dot{J}_{GW} = -\frac{32}{5} \frac{G^{7/2}}{c^5} \frac{q^2}{(1+q)^4} \frac{M^{9/2}}{a^{7/2}}$$

$$\Rightarrow q \leq \frac{5}{6} + \frac{\alpha}{2}$$

$$\alpha = \frac{d \ln R}{d \ln M_{NS}}$$

$$\alpha = \frac{d \ln R}{d \ln M_{SRM}}$$

- Condition for stable mass transfer

$$q \leq \frac{5}{6} + \frac{\alpha}{2} \leftarrow \frac{d \ln R}{d \ln M_2}$$

$\frac{M_2}{M_1}$

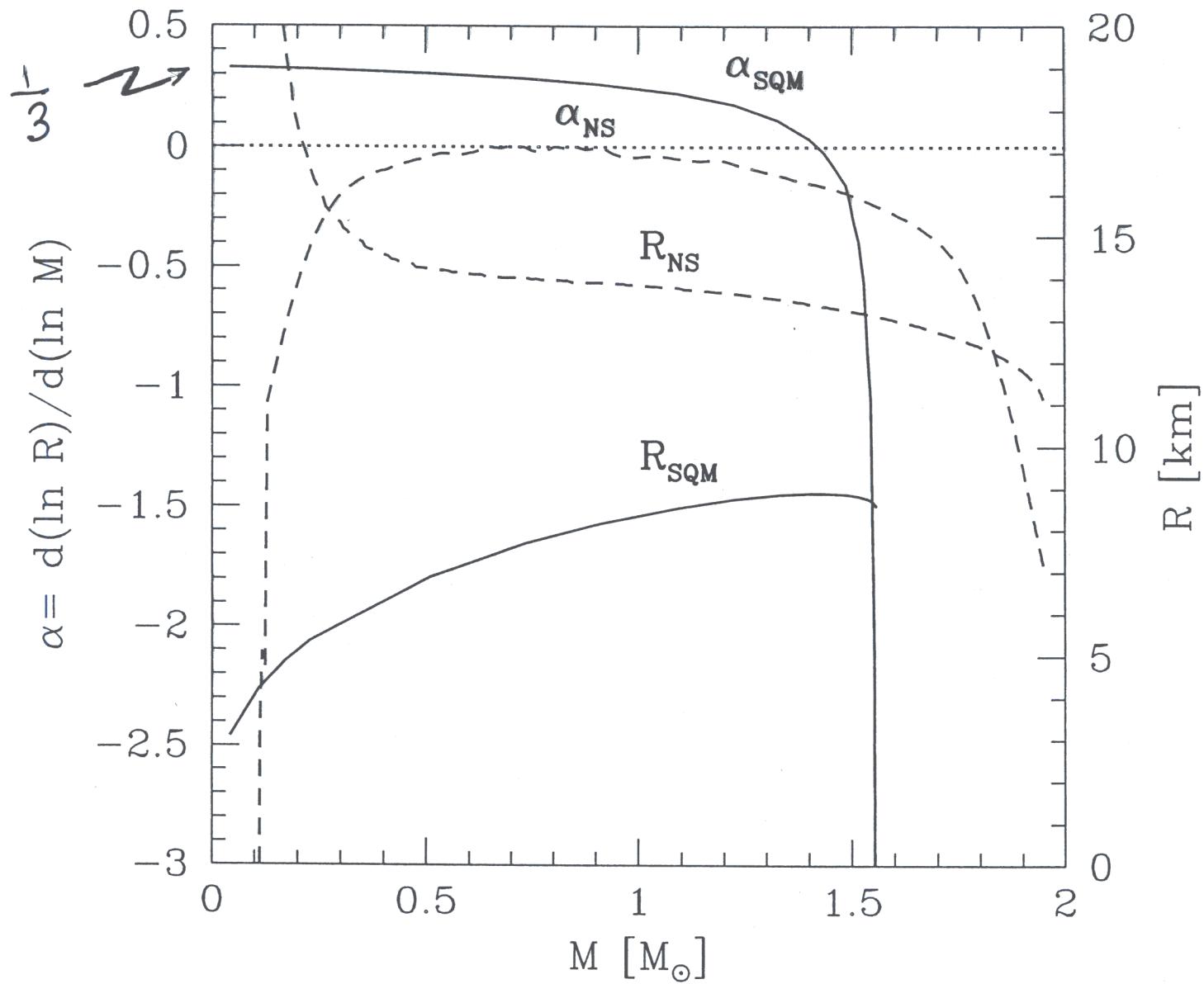
Normal:  $\alpha \approx 0 \Rightarrow q \leq \frac{5}{6} \sim 0.83$

Strange:  $\alpha \approx \frac{1}{3} - \beta \Rightarrow q \leq 1 - \beta/2$

"Basic conclusion"

SQM's allow stable mass transfer  
for a larger range of  $q$   
than do normal stars!

- For moderate mass Normal stars  
 $\alpha \leq 0$



For small to moderate mass self-bound stars (SQMs),  $\alpha \approx \frac{1}{3} - \beta$

# Test particle potentials

Newton

Paczynsky-Wiita

$$\Phi(r) = -\frac{GM}{r}$$

$$-\frac{GM}{r-r_s}$$

$$r_s = 2GM/c^2$$

$$\omega^2 = \frac{M_1 + M_2}{a^3}$$

$$\frac{M_1 + M_2}{a(a-r_G)^2}$$

$$r_G = 2(M_1 + M_2)$$

(G = c = 1)

$$\Psi(x, y) = \Phi / [(M_1 + M_2)/a]$$

$$= \frac{x_2}{\sqrt{(x+x_1^2)^2 + y^2}} - x_2 z$$

$$+ \frac{x_1}{\sqrt{(x-x_1)^2 + y^2}} - x_1 z$$

$$+ \frac{1}{2} \frac{x^2 + y^2}{(1-z)^2}$$

$$\left. \begin{array}{l} x = r_x/a \\ y = r_y/a \\ x_{1,2} = \frac{M_{2,1}}{M_2 + M_1} \\ z = r_G/a \end{array} \right\}$$

# Equipotential contours

Roche lobe

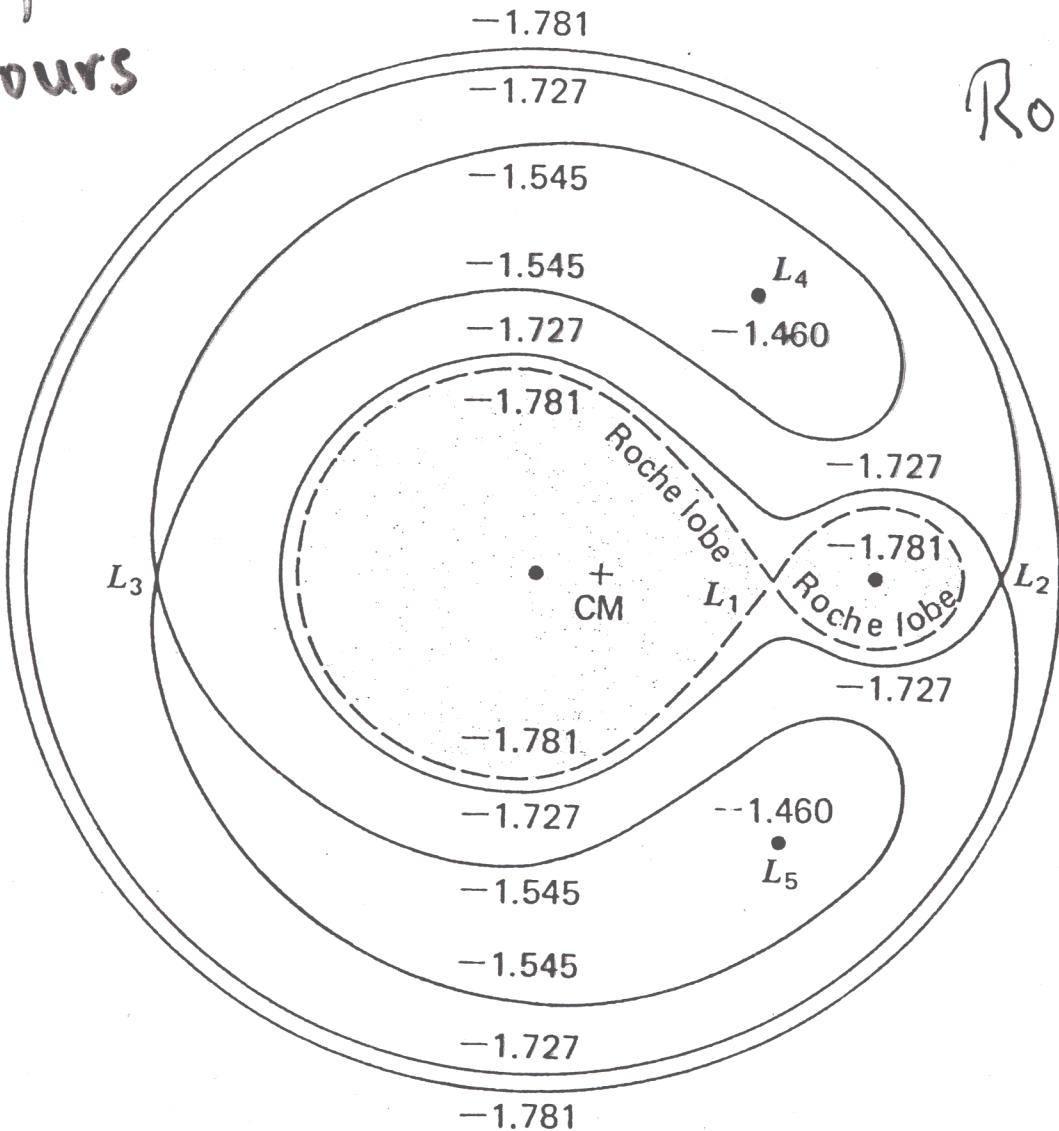
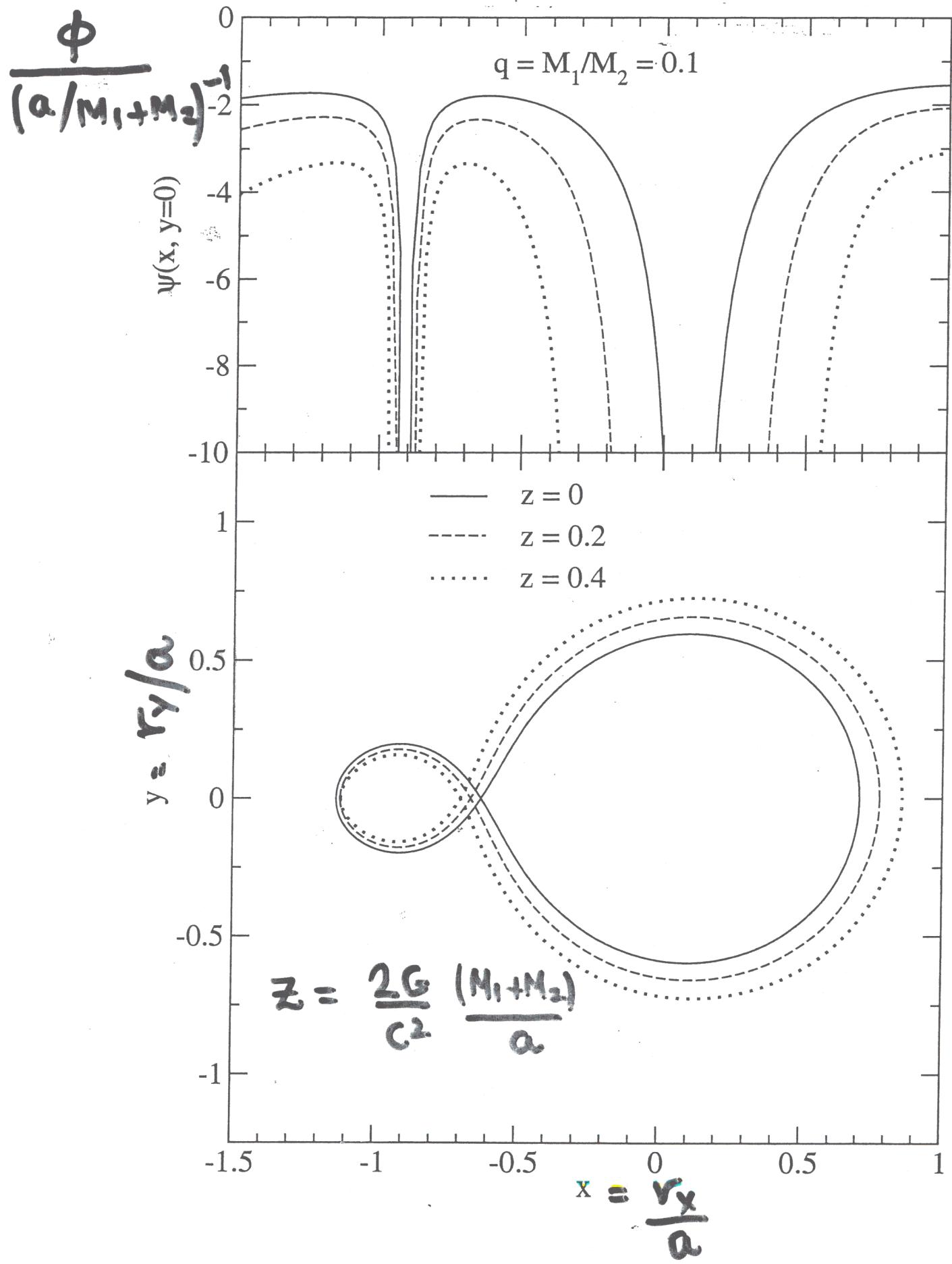
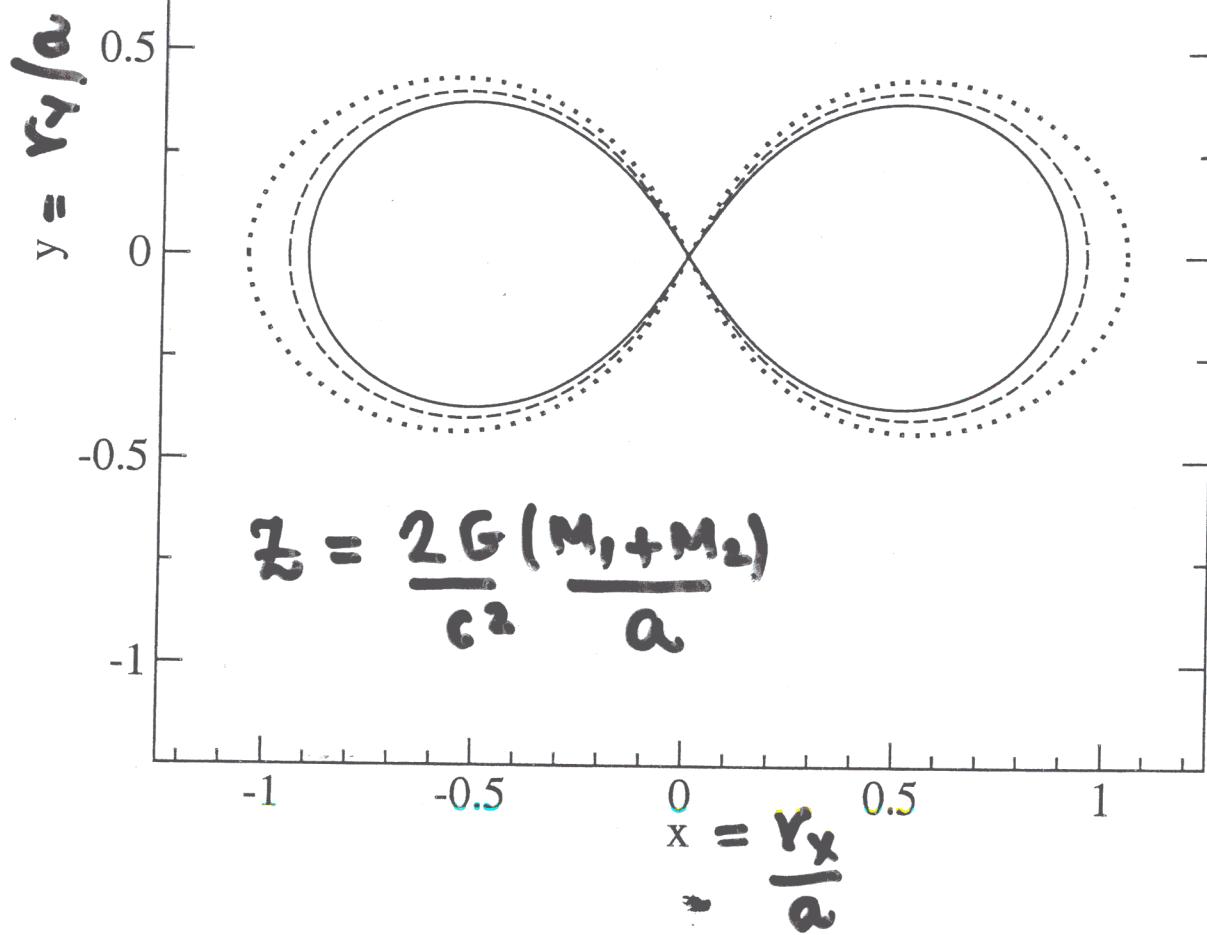
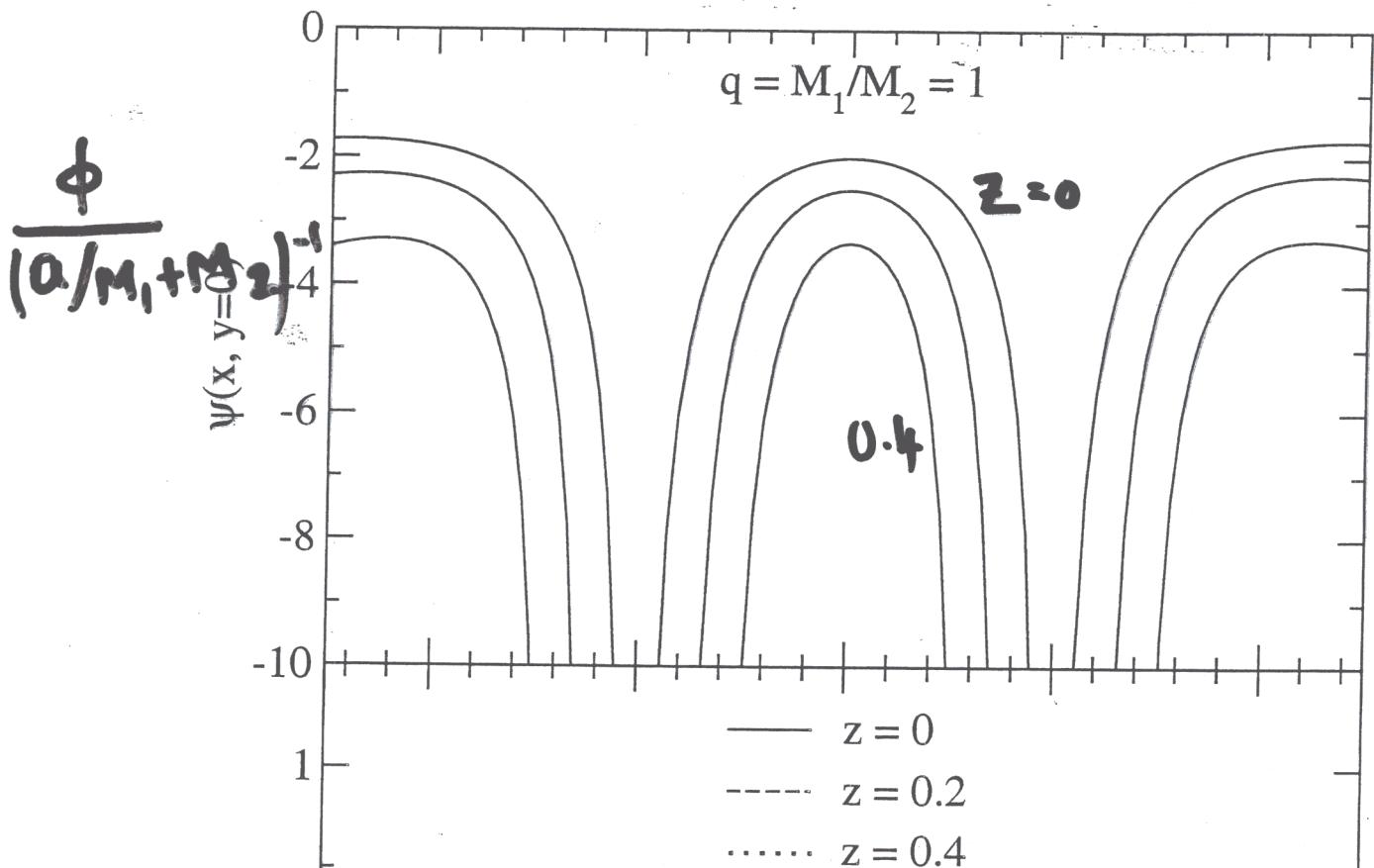


Figure 13.12 The equipotentials  $\phi_{gc} = \text{constant}$  for the Newtonian gravitational + centrifugal potential in the orbital plane of the binary star system with a circular orbit. For the case shown here we have a mass ratio  $M_N : M_c = 10 : 1$ , where the normal star  $M_N$  is on the left and the compact star  $M_c$  is on the right. The equipotentials are labeled by their values of  $\phi_{gc}$  measured in units of  $(M_N + M_c)/a$ , where  $a$  is the separation of the centers of mass of the two stars. The innermost potential shown is the "Roche lobe" of each star. Inside each Roche lobe, but outside the stellar surfaces, the potential  $\phi_{gc}$  is dominated by the "Coulomb" ( $1/r$ ) field of the star, so the equipotential surfaces are spheres. The potential  $\phi_{gc}$  has local stationary points ( $\nabla\phi_{gc} = 0$ ), called "Lagrange point" locations marked  $L_j$ . [From Novikov and Thorne (1973).]





- Roche lobe radius:

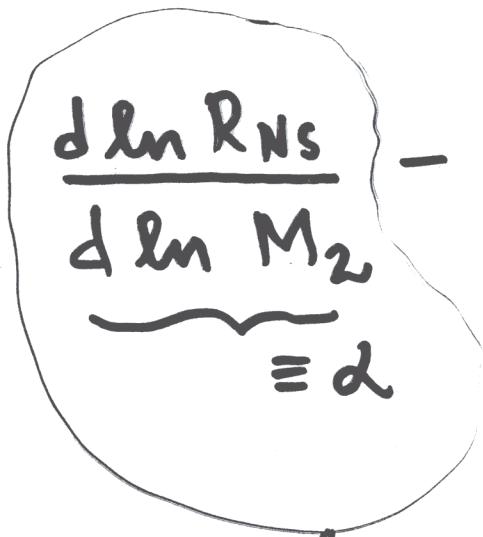
$$R_L = 0.46 a \left( \frac{q}{1+q} \right)^{1/3}$$

(Paczynski)

- Set  $R_L = R_{NS}$  or SQM

at the moment of mass transfer

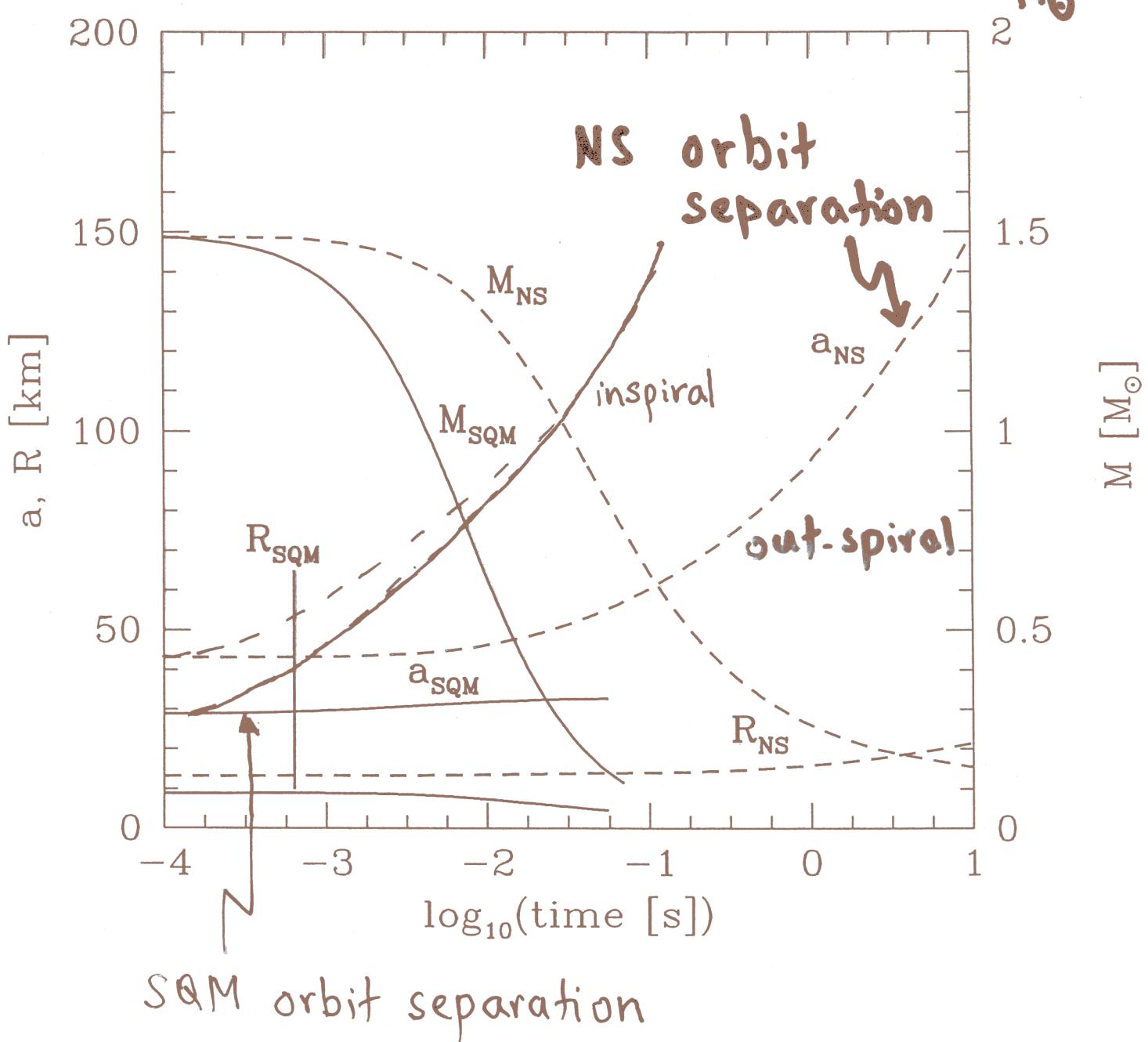
$$\frac{d \ln a}{d \ln M_2} = \underbrace{\frac{d \ln R_{NS}}{d \ln M_2}}_{\equiv d} - \frac{1}{3} = \left( d - \frac{1}{3} \right)$$



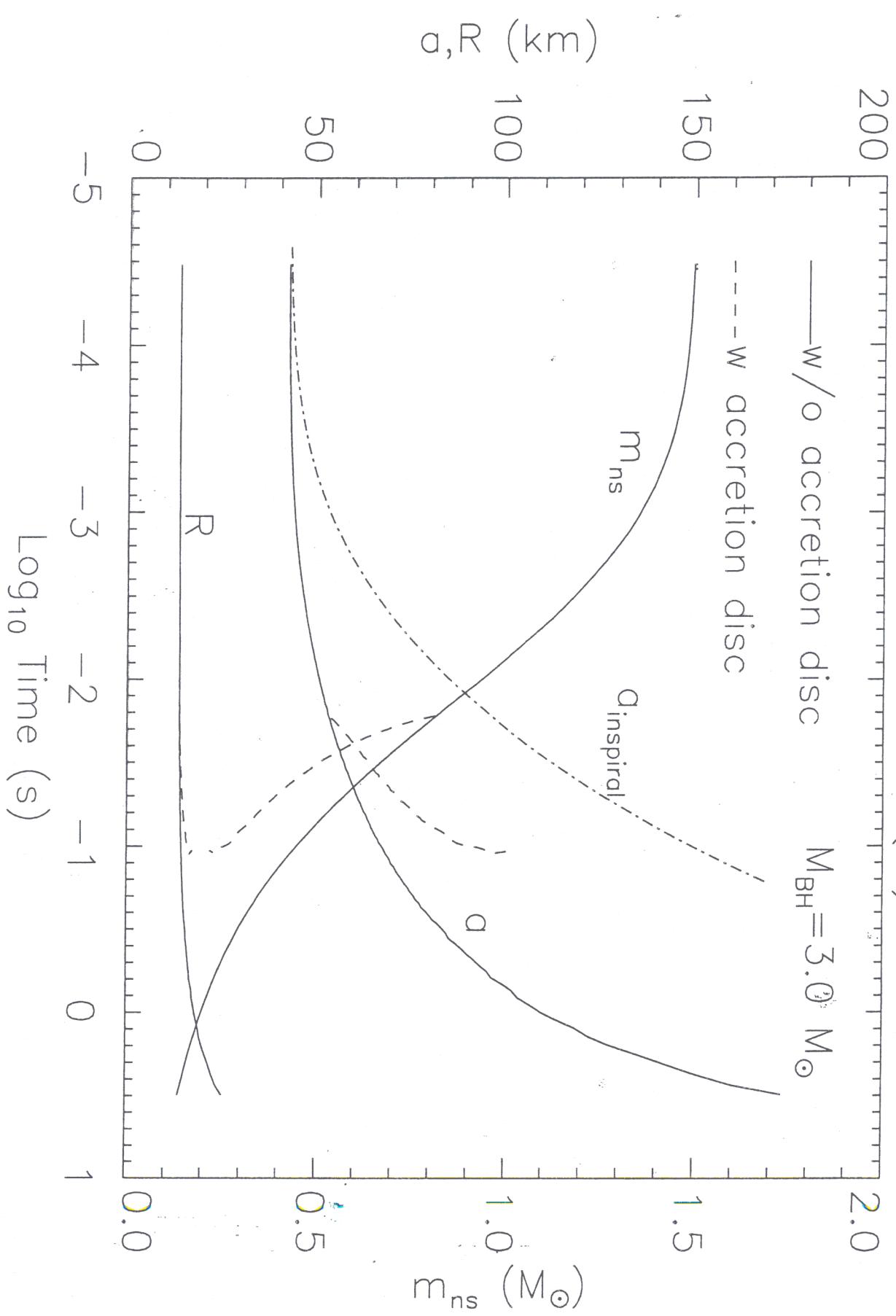
Connection  
to the EOS

- Evolution of binaries intent on merging

$$\frac{M_{BH}}{M_\odot} = 3.5$$

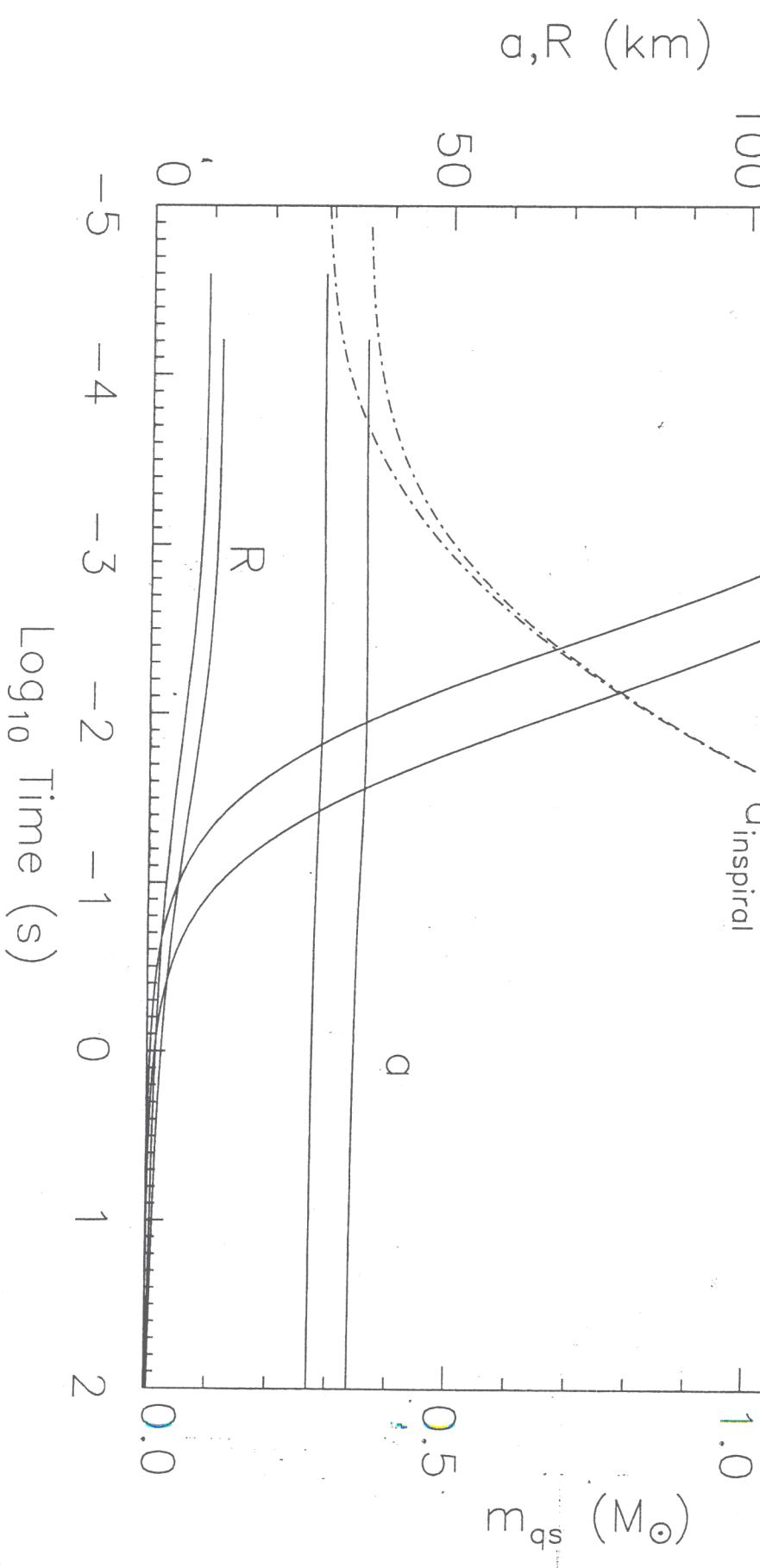


# Neutron Star R(M)

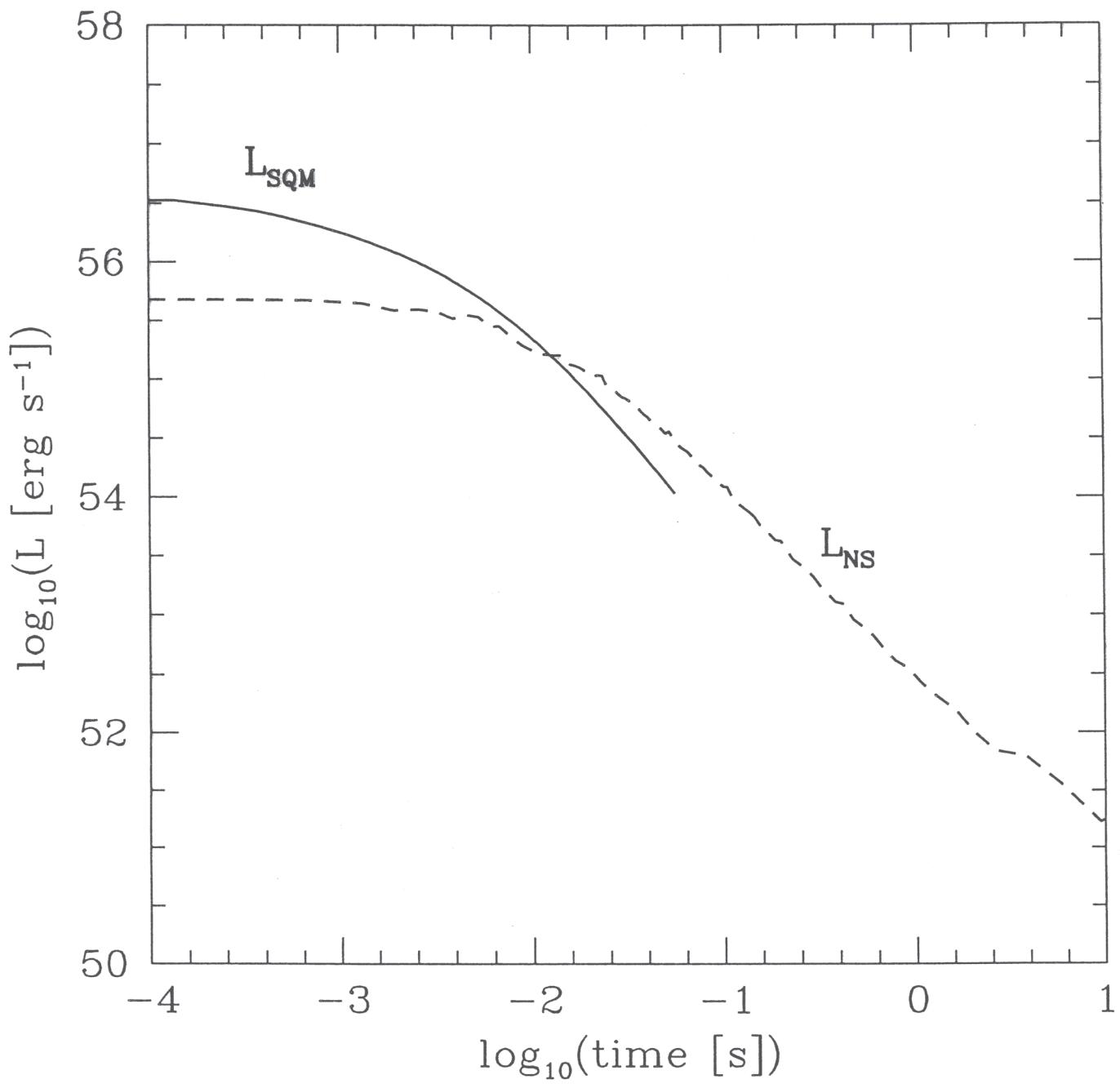


# Quark Star R(M)

$M_{BH} = 3.0 M_{\odot}$



# Luminosities in Gravitational Radiation



$$\tau_{\text{SQMs}} \ll \tau_{\text{NSs}}$$

Relative magnitudes depend on  
( $M, R$ )